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## Small-Scale CW HF(DF) Chemical Laser

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Interim Report

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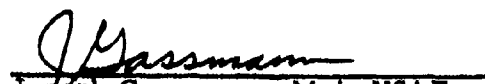
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## PREFACE

The many helpful discussions with Dr. J. Hinchey during the initial design of the laser are gratefully appreciated.

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## I. INTRODUCTION

Continuous-wave HF and DF chemical lasers have the potential for high output power; devices with output powers of the order of 1 kW have been reported.<sup>1</sup> However, there is a need for a small-scale (i.e., output power on the order of 1 W) amplitude- and frequency-stable cw chemical laser for a variety of applications. These applications include such chemical laser diagnostic studies as small-signal-gain measurements,<sup>2</sup> atmospheric propagation, and window-material and optic-coating development. A stable HF(DF) laser is also needed for use as the local oscillator in a laser-sun heterodyne radiometer detection array employed in high-resolution atmospheric-transmission measurements.<sup>3</sup>

The use of small-scale chemical lasers for diagnostic studies was reported by Cool<sup>4</sup> and Hinchey.<sup>5</sup> Cool used a microwave discharge to generate F atoms, whereas Hinchey used a dc discharge. The construction and performance of a small-scale cw HF(DF) laser that is similar to the device described by Hinchey is described here.

## II. LASER DESIGN

The laser is shown in Figs. 1 and 2. The device can be divided conveniently into the five regions (i.e., the discharge region, inlet region, lasing zone, heat exchanger, and exhaust). The discharge region provides an F-atom flow for the chemical reaction by means of a high-voltage discharge in a water-cooled pyrex tube containing  $\text{SF}_6$ , He, and  $\text{O}_2$  that generates vibrationally excited HF or DF molecules. The He serves as a thermal diluent (and also reduces the required discharge voltage), whereas the  $\text{SF}_6$  is the F-atom source. Oxygen is added to the flow to increase the F-atom concentration and to eliminate S deposition through exothermic reaction with the S atoms formed in the discharge. The mixed gases are injected into one end of the discharge tube upstream of a group of eight symmetrically deployed 1/16-in. -diameter by 6-in.-long Ni wires that serve as discharge anodes. The operation of the Ni wires as cathodes resulted in metal loss and reduced power output. The discharge uniformity depends heavily upon the flow uniformity of the gases at the electrode tips; therefore, the gas injection into the discharge tube is made as uniform as possible. The long anodes are gas-cooled in the region upstream of the discharge. The Ni anodes are individually connected through ballast resistors (50 k $\Omega$ , 100 k $\Omega$ , or 200 k $\Omega$ ) to the positive terminal of the power supply. The ballast resistor values can be varied to accommodate power-supply characteristics and laser operating conditions. The ballast resistor bank is forced-air-cooled.

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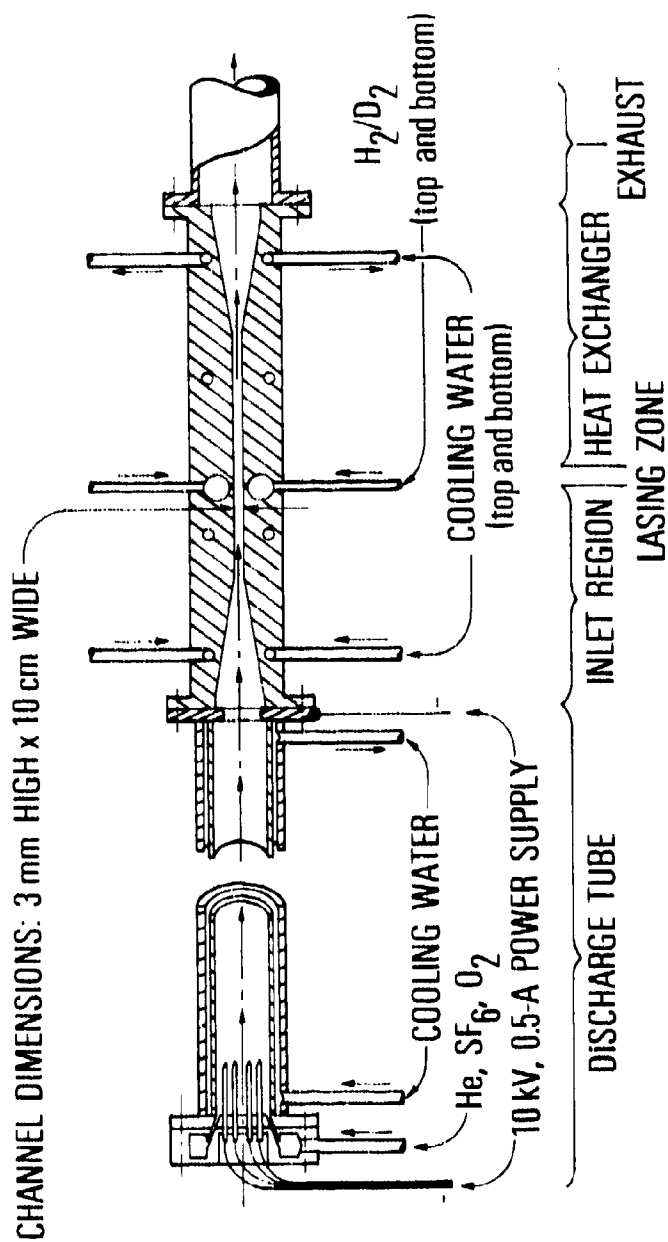


Figure 1. Stable HF(DF) Chemical Laser

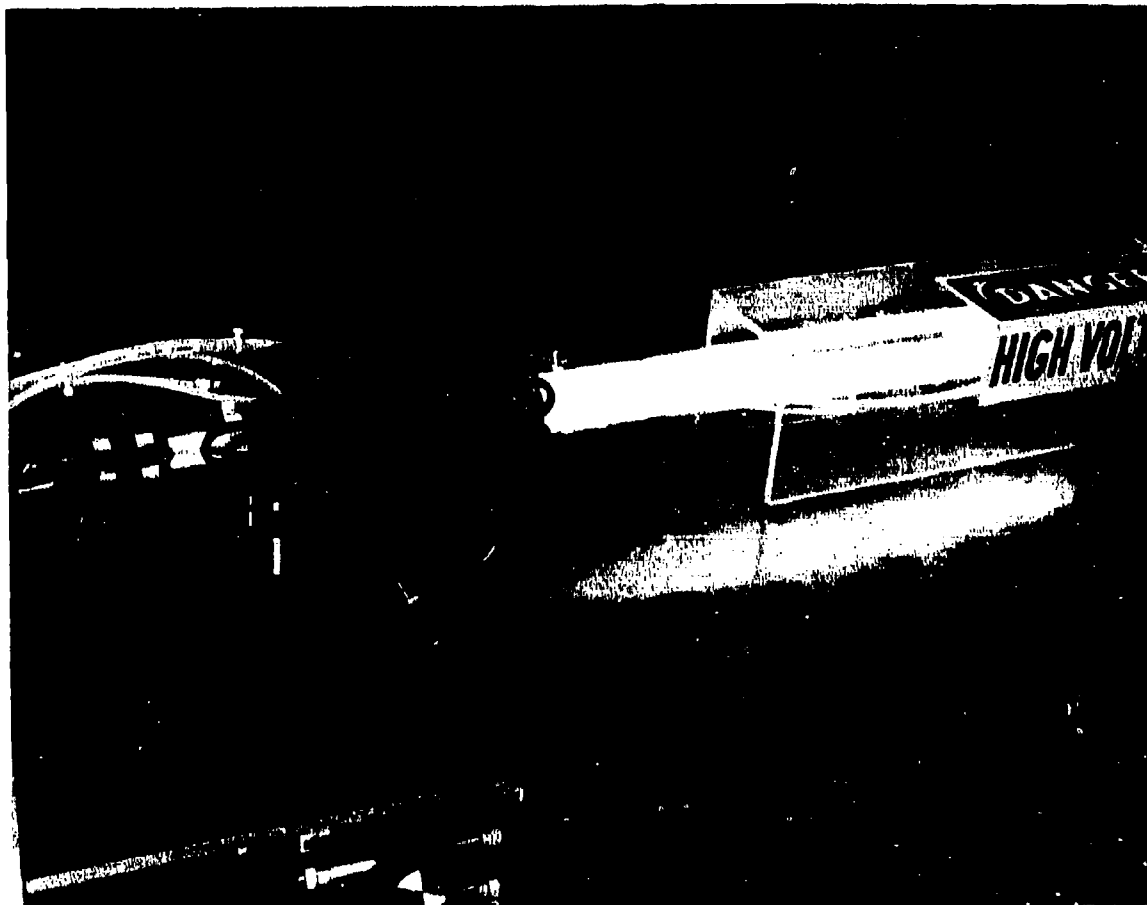


Figure 2. Laser Discharge Tube and Cavity

The discharge takes place along a central core in the pyrex tube and terminates on a water-cooled Cu hole electrode (3/4-in. diameter) at the entrance to the laser inlet region. The pyrex tube has a 1-in. i.d., 0.062-in. wall thickness and nominally is used in two lengths, 12 and 24 in., with 10- and 20-kV power supplies, respectively. An  $\text{Al}_2\text{O}_3$  cylindrical shell is inserted inside the pyrex tube at the junction with the cathode to eliminate discharge attachment at the pyrex-Cu cathode interface. The water jacket surrounding the pyrex tube is made of lucite.

The hot flow containing F atoms passes through the cathode hole into the inlet region where the flow cross section is changed from a square geometry of the same nominal cross section as the discharge tube to the channel cross section 10 by 0.3 cm. The  $\text{H}_2$  is injected along a line transverse to the flow at both the top and bottom of the channel through 40 evenly spaced 0.0135-in.-diameter holes in each wall. The  $\text{H}_2$  is fed from plenums immediately above and below the injection hole array. The lasing zone is immediately downstream of the  $\text{H}_2$  injection ports, where the mixing of the  $\text{H}_2$  and F takes place. The remaining portion of the channel downstream of this region and the transition section in which the rectangular channel geometry is returned to a circular geometry for convenience in ducting serve as a heat exchanger to reduce the temperature of the heated gas to  $\sim 200^\circ\text{C}$ . Further temperature reduction of the gas is accomplished in a downstream heat exchanger.

The inlet region, the mixing and lasing zone, and the heat exchanger are fabricated in a single Amsulf Cu block 8 by 4 by 1-1/2 in. The assembly is water-cooled and contains the  $\text{H}_2$  plenums. The laser is exhausted to a 150

cfm) vacuum pump through a conventional 1-1/2 in. i.d. flexible vacuum hose. A surge tank in the exhaust is unnecessary for damping pump pressure fluctuations since no pressure variations at the laser are detectable with the undamped system operating at normal flow rates. Brewster-angle  $\text{BaF}_2$  windows are mounted on channel sidewalls that are O-ring sealed to the channel body. The windows are centered along the gain region in the channel, nominally, 1-mm downstream of the  $\text{H}_2$  injection holes. Laser-beam polarization can be controlled by the rotation of the Brewster-angle windows for proper orientation with respect to the gratings. In an early experiment, a small purge flow of He was injected at the Brewster windows to reduce the accumulation of ground-state HF or DF, but the use of the flow did not affect laser performance and, hence, was discontinued. The optical cavity is constructed from Burleigh optical erector set components and consists of stainless-steel star gimbal mounts on stainless-steel rod plates that are clamped to four 1-in. -diameter Invar rods. The laser Cu body is suspended between the optical cavity rods with the gain region colinear with the optical-cavity centerline. Multiline power can be coupled from the laser by means of a partially transmitting mirror used in conjunction with a totally reflecting mirror, one of which is a spherical mirror of radius 1 to 10 m and the other, an optical flat. Single-line operation can be accomplished by use of a 300- or 600-lines/mm grating in the autocollimation position with zero-order coupling for lasing at the desired wavelength and with a spherical totally reflecting mirror used as the other cavity component.

### III. OPERATING CHARACTERISTICS

The laser is generally used with a 150-cfm vacuum pump and a 20-kV 500-mA power supply. A 24-in. long discharge tube is employed with power supplies of 10 kV or greater terminal voltage. A 12-in.-long discharge tube is used with 5 and 10 kV open-circuit-voltage power supplies capable of operating at  $\sim 1$  A current levels. Maximum power is obtained with the higher voltage operation, however. Two 100-k $\Omega$ , 250-W ballast resistors are connected in series with each of the eight Ni wire electrodes. Thus, the total ballast resistance in series with the discharge is 25 k $\Omega$ . This ballast resistance value yields optimum electrical efficiency when used with a 20-kV open-circuit-voltage power supply (50 k $\Omega$ /leg is optimum with a 10-kV supply). The multiline HF laser performance for several total-ballast resistance values ranging from 12.5 to 37.5 k $\Omega$  and including the electrical-power distribution in the ballast-discharge circuit, the electrical to laser-power conversion efficiency, and the channel operating pressures 1/2-in. upstream of the laser cavity and at the channel exit is given in Table I. Maximum laser power for the standard HF operating condition (Table II) is 12 W. The maximum overall electrical-to-laser power conversion efficiency is obtained with the 25-k $\Omega$  ballast resistor and is equal to 1/2%. The power distribution between ballast and discharge is 65/44 for this ballast resistor. The power supply operates at 15.6 kV and 350 mA. Gas pressure is  $\sim 23.5$  Torr in the laser-cavity region. Gas temperatures can also be measured by means of thermocouples placed in the flow at the channel entrance and exit

Table I. Laser, Ballast Resistor, and Power Supply Operating Characteristics

Ballast Resistance ( $\Omega$ )	Current (mA)	Circuit			Power			Efficiency (%)	Pressure	
		Power Supply Terminals (kV)	Ballast Resistor (kV)	Laser Discharge (kV)	Power Supply (kW)	Ballast Resistor (kW)	Laser Discharge (kW)		Laser Output (W)	Exit Loss (W)
HF LASER <sup>a</sup>										
12.5	500	12.5	7.25	6.25	6.25	3.125	3.125	0.36	12.0	25.2 <sup>a</sup>
16.7 <sup>b</sup>	400	14.5	7.5	7.5	5.60	3.0	2.60	0.43	12.3	25.0
25.0	350	15.0	7.25	7.45	5.46	3.06	2.40	0.50	12.0	23.4
37.5	300	17.0	11.25	7.75	5.76	3.375	2.325	0.43 <sup>b</sup>	10.0 <sup>b</sup>	23.9
DF LASER <sup>c</sup>										
12.5	500	12.7	6.30	6.40	6.35	3.15	3.20	0.32	7.0	25.0
16.7 <sup>b</sup>	400	14.7	7.50	7.20	5.65	3.00	2.60	0.24	7.0	24.5
25.0	350	16.4	7.25	7.65	5.74	3.06	2.68	0.26	7.0	24.0
37.5	300	19.4	7.75	7.65	5.94	2.53	2.51	0.20 <sup>c</sup>	5.0 <sup>c</sup>	22.3

<sup>a</sup> Laser operating conditions: 24-in.-long discharge tube; laser cavity: 40-in.-radius-of-curvature of total reflector and 157-coupling mirror;  $\dot{m}_{He} = 0.654$ ;  $\dot{m}_{D_2} = 0.030$ ;  $\dot{m}_{O_2} = 0.160$  mSg; power supply: 20 kV open-circuit voltage.

<sup>b</sup> Laser output power not optimized because of power supply voltage limitation (20 kV).

<sup>c</sup> Laser operating conditions: 24-in.-long discharge tube; laser cavity: 40-in.-radius-of-curvature of total reflector and 1% coupling mirror;  $\dot{m}_{He} = 0.9405$ ;  $\dot{m}_{D_2} = 0.044$ ;  $\dot{m}_{O_2} = 0.150$  mSg; power supply: 20 kV open-circuit voltage.

<sup>a</sup> Laser operating conditions: 24-in.-long discharge tube; laser cavity: 40-in.-radius-of-curvature of total reflector and 157 coupling mirror.

<sup>b</sup>  $m_{He} = 0.044$ ;  $m_{H_2} = 0.030$ ;  $m_{O_2} = 0.160$ ;  $m_{SF_6} = 0.070$  g/sec; power supply: 20 kV open circuit voltage.

<sup>c</sup> Laser output power not optimized because of power supply voltage limitation (20 kV).

<sup>d</sup> Laser operating conditions: 24-in.-long discharge tube; laser cavity: 40-in.-radius-of-curvature of total reflector and 7% coupling mirror.

<sup>e</sup>  $m_{He} = 0.0405$ ;  $m_{D_2} = 0.044$ ;  $m_{O_2} = 0.150$ ;  $m_{SF_6} = 0.050$  g/sec; power supply: 20 kV open circuit voltage.

Table II. Standard HF Laser Operating Conditions

Electrical Power	Voltage = 15.6 kV, current = 350 mA, $R_D = 25 \text{ k}\Omega$ , 24-in. -discharge tube length
Optical Cavity	40-in. -radius-of-curvature of total reflector and 15% dielectric output coupler separated by 16 in.
Cavity Gas Conditions	Pressure: 1/2 in. upstream of lasing region = 23.5 Torr, channel exit = 15.7 Torr. Temperature: channel entrance = $\sim 1000^\circ \text{K}$ , lasing region = $\sim 700^\circ \text{K}$ , channel exit = $\sim 500^\circ \text{K}$ .
Gas Mass Flow	$\dot{m}_{\text{O}_2} = 0.180 \text{ g/sec}$ $\dot{m}_{\text{He}} = 0.054 \text{ g/sec}$ $\dot{m}_{\text{H}_2} = 0.026 \text{ g/sec}$ $\dot{m}_{\text{SF}_6} = 0.870 \text{ g/sec}$

and spectrographically in the lasing zone from HF laser line width measurements obtained by the piezoelectric controlled variation of the cavity length. Channel-entrance temperatures are typically  $\sim 1000^\circ\text{K}$ , whereas the lasing zone temperature is  $\sim 700^\circ\text{K}$ . This relatively high laser reaction temperature results in a peaking of the laser lines at rotational quantum number  $j = 7$  in HF operation and  $j = 10$  in DF operation.

Similar data for DF laser operation are given in Tables I and III. Maximum laser power for the standard DF operating conditions is 7 W. Again, the maximum overall electrical-to-laser power conversion efficiency is obtained with the 25-k $\Omega$  ballast resistor and is equal to 0.26%. Power-supply voltage is 16.4 kV, and discharge current is again 350 mA. Gas pressure is  $\sim 24$  Torr and gas temperature  $\sim 700^\circ\text{K}$ .

Maximum multiline laser output power ever observed in this laser is 15 W with  $\text{SF}_6$  used as the F-atom source. Substitution of  $\text{NF}_3$  for  $\text{SF}_6$  results in a 18-W HF and DF laser output. Several attempts were made to increase laser power by redesign, which included increasing the number of the injection holes from 40 to 120 per side, but no net gain in laser power was observed. The laser performance appears to be F-atom-limited and not diffusion-limited. The reduction of performance with mirror, window, or flow degradation is always much greater with the DF laser than the HF because of the lower gain per line in the DF laser.

The static-pressure measurements obtained at entrance and exit points on the channel as well as at the pump entrance flange are given in Fig. 3. Pressure measurements are made 2.5 and 3.0 in. downstream of the



Table III. Standard DF Laser Operating Conditions

Electrical Power:	Voltage = 16.4 kV, current = 350 mA, $R_B = 25 \text{ k}\Omega$ , 24-in. -discharge tube length
Optical Cavity:	40-in. -radius-of-curvature of total reflector and 7% dielectric output coupler separated by 16 in.
Cavity Gas Conditions	Pressure: 1/2 in. upstream of lasing region = 24.0 Torr, channel exit = 14.0 Torr. Temperature: channel entrance = $\sim 1000^\circ\text{K}$ , lasing region = $\sim 700^\circ\text{K}$ , channel exit = $\sim 500^\circ\text{K}$ .
Gas Mass Flows	$\dot{m}_{\text{O}_2} = 0.150 \text{ g/sec}$ $\dot{m}_{\text{He}} = 0.0405 \text{ g/sec}$ $\dot{m}_{\text{D}_2} = 0.044 \text{ g/sec}$ $\dot{m}_{\text{SF}_6} = 0.980 \text{ g/sec}$

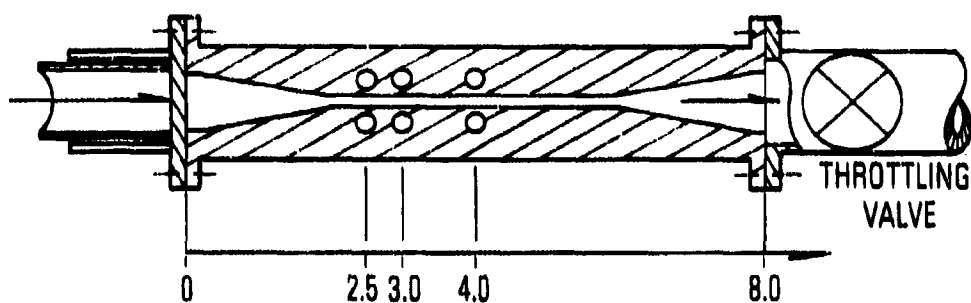
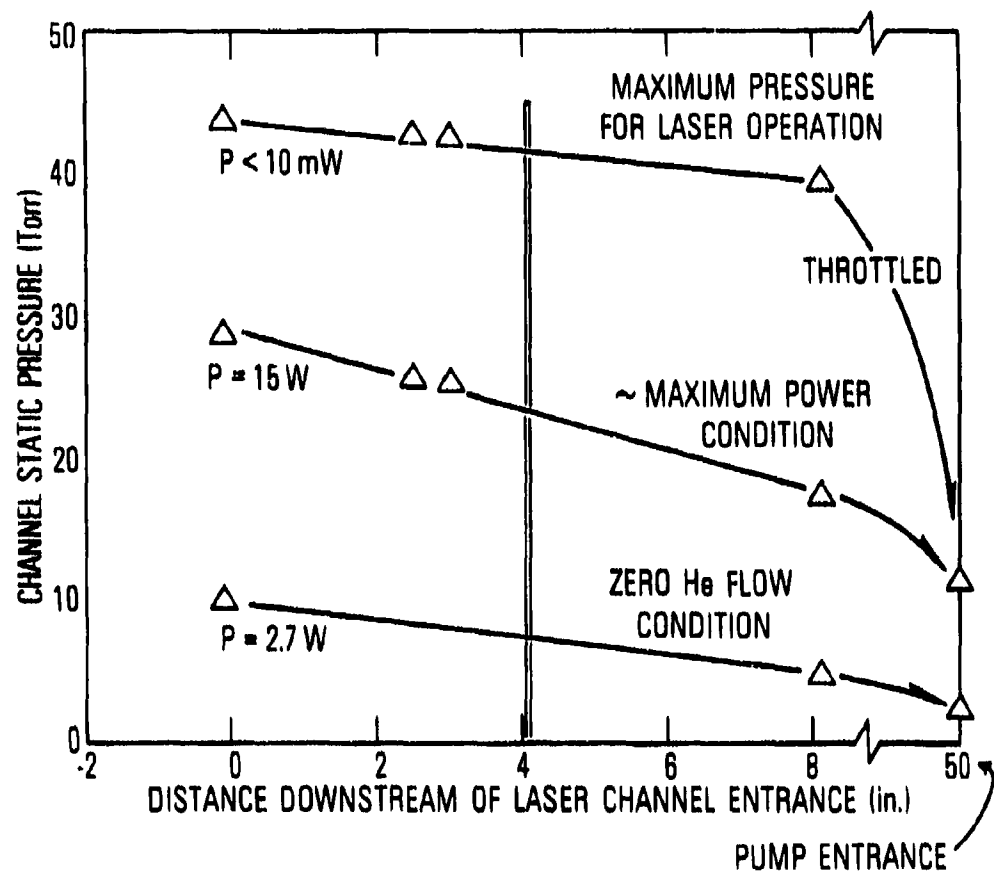


Figure 3. Static Pressure in Channel for Various Operating Conditions

channel entrance in a channel configuration (cross section, Fig. 3). This configuration is identical to the standard configuration (Fig. 1) except for the placement of three sets of  $H_2$ -injection plenums and the associated 40 injection holes at each plenum. It is used to determine the dependence of the laser-power output on the longitudinal injection position of the  $H_2$  in the channel. However, no significant power difference is observed in this configuration at any of the three locations. Also included in Fig. 3 are the static-pressure measurements for a zero He flow run condition. Multiline HF power is measured at 2.7 W for this situation. In addition, the maximum allowed pressure for laser operation is determined by the downstream throttling of the flow under maximum power flow conditions. Laser cavity pressure is  $\sim 41$  Torr at a power level of  $< 10$  mW for this condition.

Normalized output power curves for the various gases used in HF-laser operation are shown in Figs. 4 through 7. The intercepts of the curves with the ordinate axes indicate the criticality of the various gases in the lasing process. For example, both  $O_2$  and He curves exhibit lasing at 16% full power at zero He or  $O_2$  flow, indicating the roles of these gases as auxiliaries to the fuel and oxidizer. However, the sixfold increase in laser power resulting from the addition of each of these gases to the flow illustrates the importance of the influence of these two gases on F-atom generation or thermal control, or both. The  $H_2$  intercept is at zero, as expected for molecular fuel injection. On the other hand, the  $SF_6$  intercept is negative, indicating a recombination loss of F atoms to below threshold requirements in the channel prior to fuel injection.

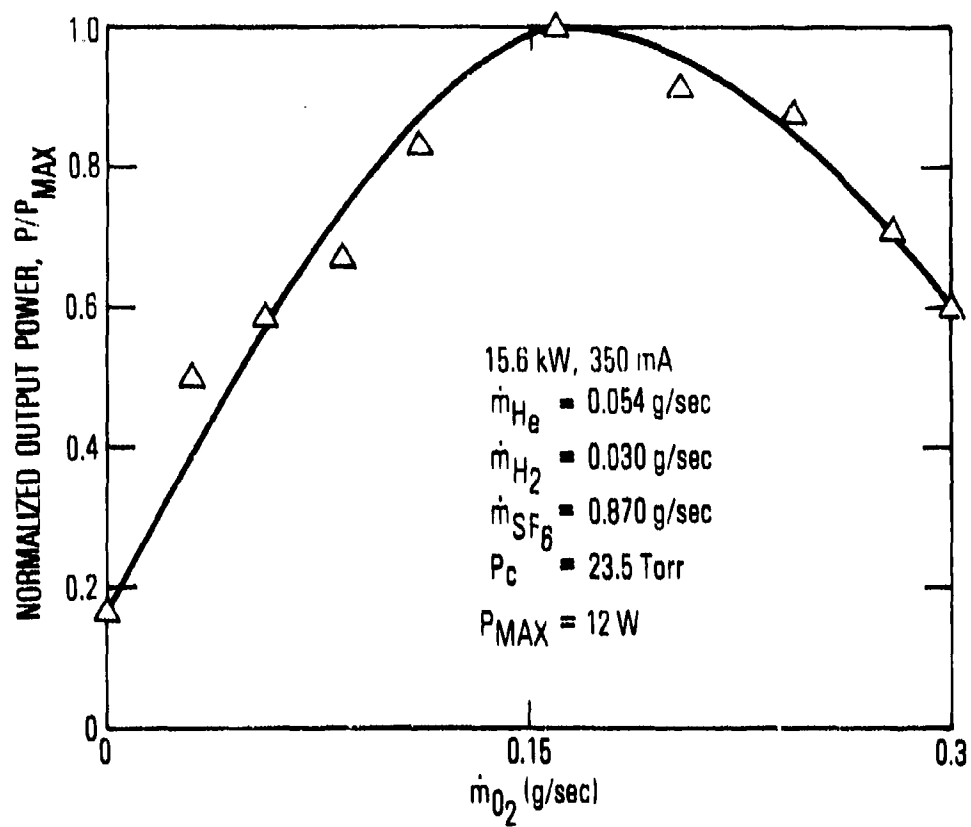


Figure 4. Variation of HF Laser Power with  $\text{O}_2$  Mass Flow

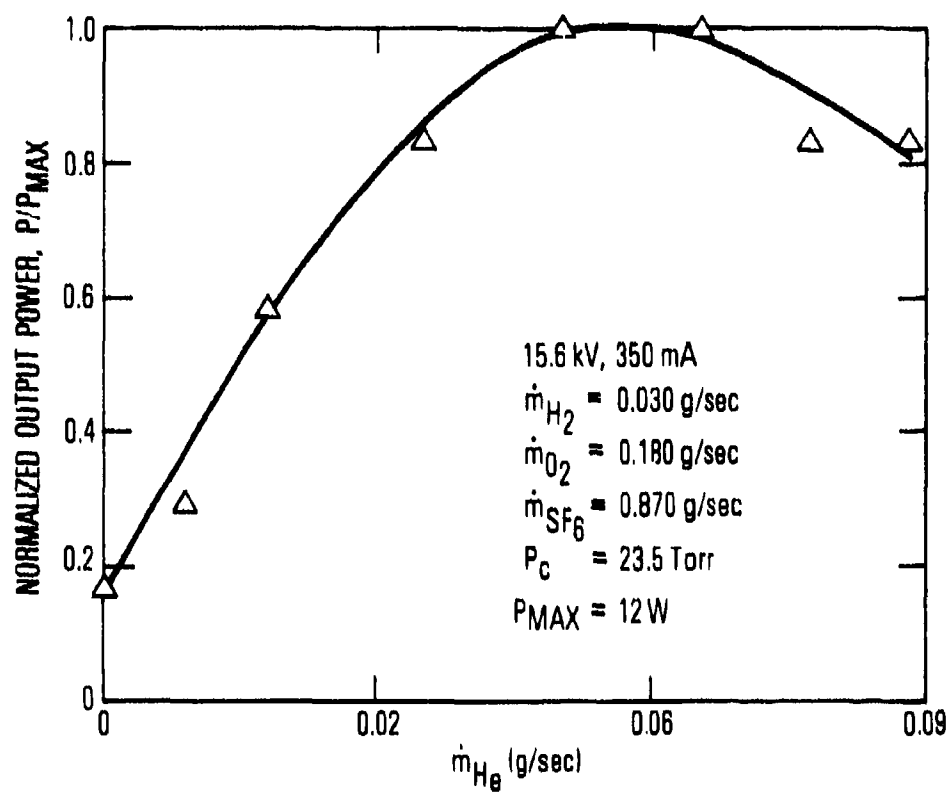


Figure 5. Variation of HF Laser Power with He Mass Flow

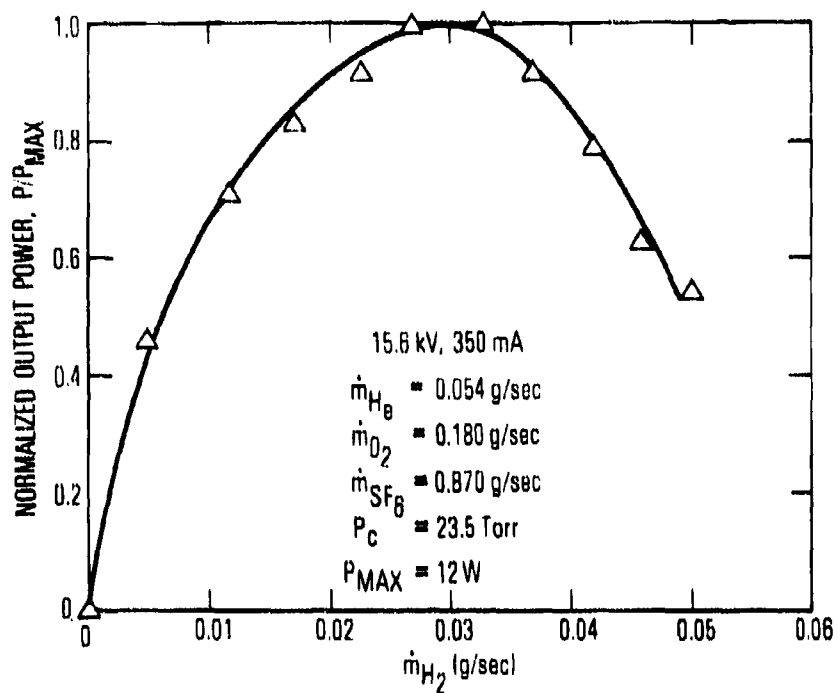


Figure 6. Variation of HF Laser Power with  $\text{H}_2$  Mass Flow

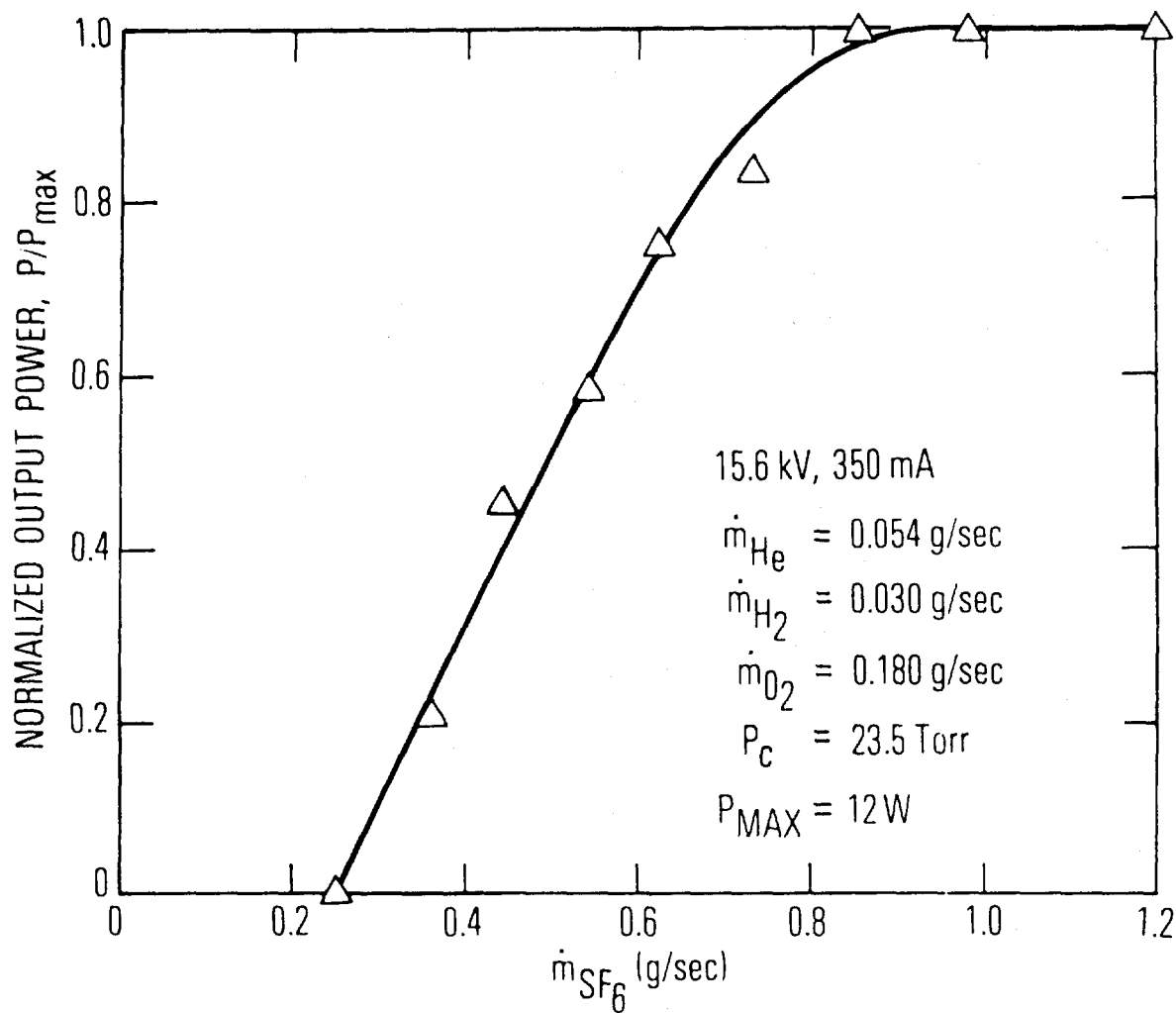


Figure 7. Variation of HF Laser Power with  $\text{SF}_6$  Mass Flow

Figures 8, 9, 10, and 11 are normalized laser output curves for the various gases in DF laser operation. Only the  $\text{SF}_6$  curve is identical to the HF counterpart. Comparison of the He curves for both HF and DF indicates a 16 versus 50% full-power zero-flow intercept. The  $\text{O}_2$  curves for HF and DF indicate an even more dramatic difference in zero  $\text{O}_2$  flow power (i. e., 16 versus 66%). Thus, the effect of He and  $\text{O}_2$  is less pronounced for the DF laser than for the HF laser. Furthermore, the  $\text{D}_2$  curve peaks at 0.7 times the molar flow of  $\text{H}_2$ . Detailed consideration of these differences in gas-flow power behavior between HF and DF lasers will be further studied.



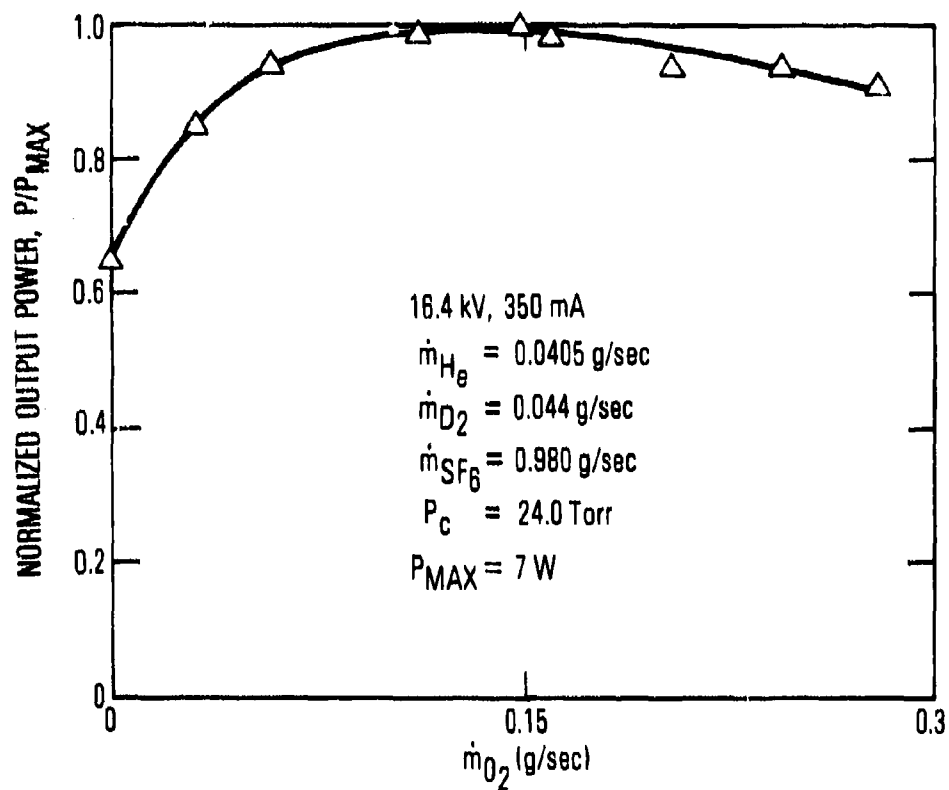


Figure 8. Variation of DF Laser Power with  $O_2$  Mass Flow

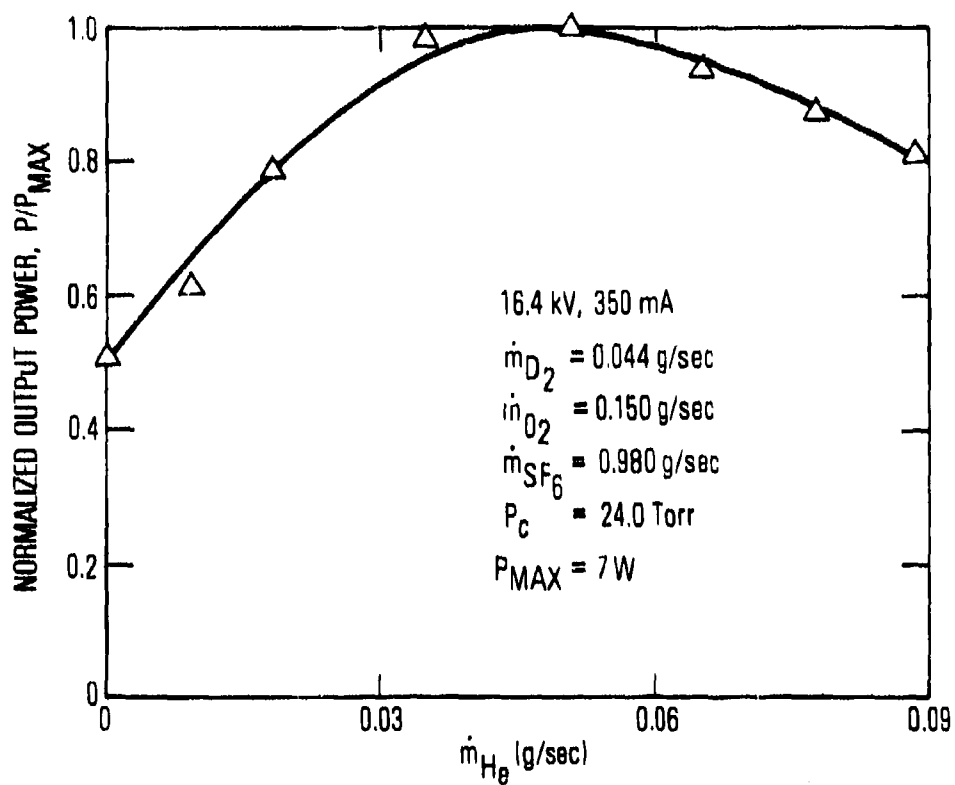


Figure 9. Variation of DF Laser Power with He Mass Flow

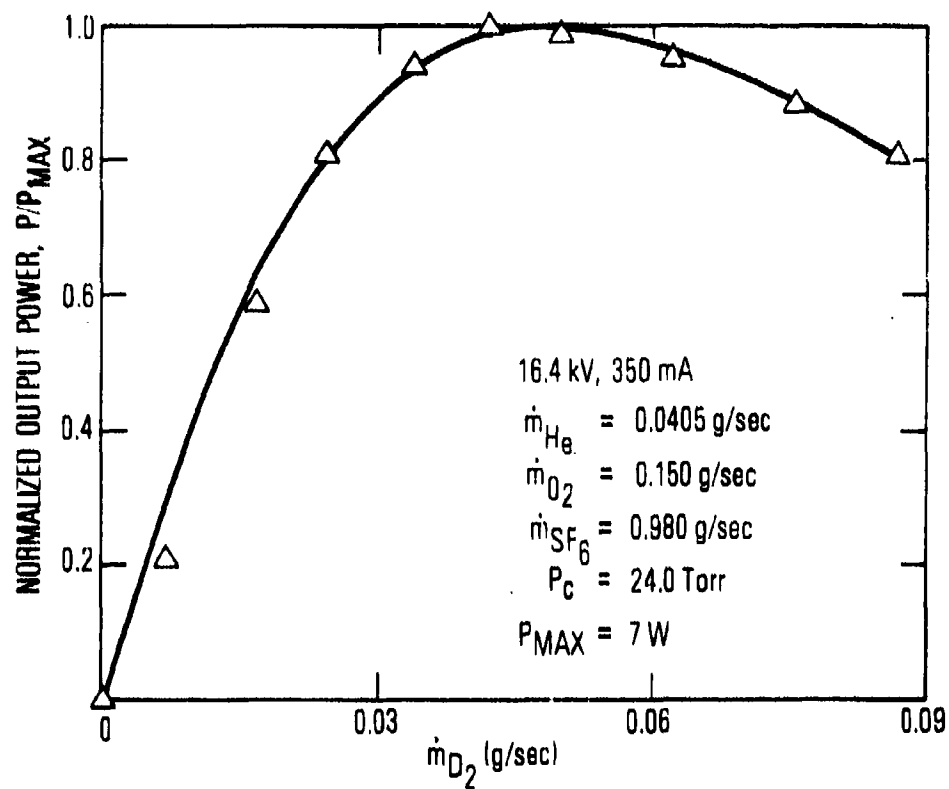


Figure 10. Variation of DF Laser Power with  $D_2$  Mass Flow

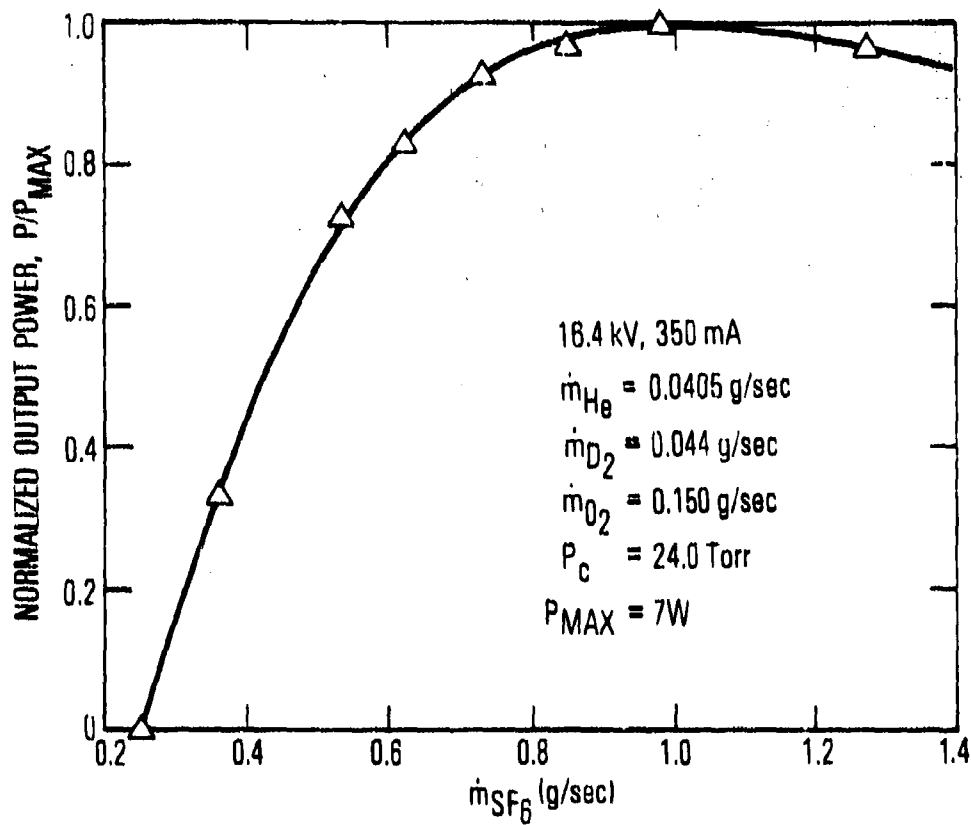


Figure 11. Variation of DF Laser Power with  $SF_6$  Mass Flow

#### IV. SINGLE-LINE PERFORMANCE

Diffraction-limited  $TEM_{00}$  single-line laser beams are obtained directly from the laser cavity. Aperture-limiting of the lasing mode to  $TEM_{00}$  is also employed but is generally not required. The laser beam is extracted as the grating zero-order beam. A PTR 92% efficient DF grating is used for power extraction for both HF and DF lines. The laser Brewster-angle windows are positioned to yield a vertically oriented E-vector that is parallel to the grating grooves for efficient HF operation. The laser lines and measured power obtained for each line at a distance 12 in. from the grating are given in Table IV. In addition, the power measurements obtained with the low-efficiency Bausch and Lomb replica grating ( $\eta \sim 80\%$ ) are tabulated. Note that the  $P_1(5)$  line is severely attenuated relative to other lines as a result of the strong atmospheric water vapor absorption. Eighteen lines are observed.

DF laser wavelengths and measured power levels obtained with the PTR 92%  $\eta$  grating are tabulated in Table V. DF lasing cannot be sustained on individual lines with the Bausch and Lomb replica grating.

Table IV. HF Laser Single-Line Performance

Line Designation P (j lower) v upper	Wavelength ( $\mu\text{m}$ )	Power PTR Grating (W)	Power Bausch and Lomb Grating (W)
P <sub>1</sub> (2)	2.5788	0.12	0.02
(3)	2.6085	0.45	0.05
(4)	2.6398	1.50	0.12
(5)	2.6727	0.15	0.15
(6)	2.7075	1.50	0.30
(7)	2.7441	2.40	0.50
(8)	2.7826	1.40	0.40
(9)	2.8231	1.00	0.30
(10)	2.8657	0.30	0.10
P <sub>2</sub> (3)	2.7275	0.20	0.05
(4)	2.7604	0.50	0.30
(5)	2.7853	1.30	0.62
(6)	2.8318	2.20	0.78
(7)	2.8706	2.20	0.95
(8)	2.9111	1.60	0.85
(9)	2.9539	0.90	0.70
(10)	2.9989	0.50	0.45
(11)	3.0461	0.30	0.20

Table V. DF Laser Single-Line Performance

Line Designation P (j lower) v upper	Wavelength ( $\mu\text{m}$ )	Power PTR Grating (W)
$P_1(5)$	3.5806	0.01
(6)	3.6128	0.025
(7)	3.6456	0.04
(8)	3.6798	0.10
(9)	3.7155	0.15
(10)	3.7520	0.20
(11)	3.7902	0.15
(12)	3.8298	0.10
(13)	3.8708	0.05
$P_2(3)$	3.6363	0.025
(4)	3.6665	0.06
(5)	3.6983	0.10
(6)	3.7310	0.15
(7)	3.7651	0.20
(8)	3.8007	0.25
(9)	3.8375	0.30
(10)	3.8757	0.32
(11)	3.9155	0.27
(12)	3.9565	0.20

Table V. DF Laser Single-Line Performance (Continued)

Line Designation P (j lower) v upper	Wavelength ( $\mu\text{m}$ )	Power PTR Grating (W)
P <sub>3</sub> (7)	3.8903	0.01
(8)	3.9272	0.02
(9)	3.9654	0.05
(10)	4.0054	0.14
(11)	4.0464	0.07



## V. LASER STABILITY

Amplitude stability is  $\sim \pm 1.5\%$  for most HF lines. The DF lines nearer threshold are more unstable, with most lines being amplitude-stable to within  $\sim \pm 5\%$ . The frequency stability of the laser is inferred initially from frequency scans of individual lines containing a Lamb dip (Fig. 12).<sup>6</sup> The clearly discernible Lamb dip feature in the trace indicates a frequency stability an order of magnitude smaller than the average line width (i.e., approximately  $\pm 10$  to  $30$  MHz). This natural-frequency stability estimate has since been confirmed by a self-beating method<sup>7</sup> to be  $30$  MHz FWHM. Line widths are typically  $\sim 400$  MHz for HF and  $\sim 300$  MHz for DF lines.

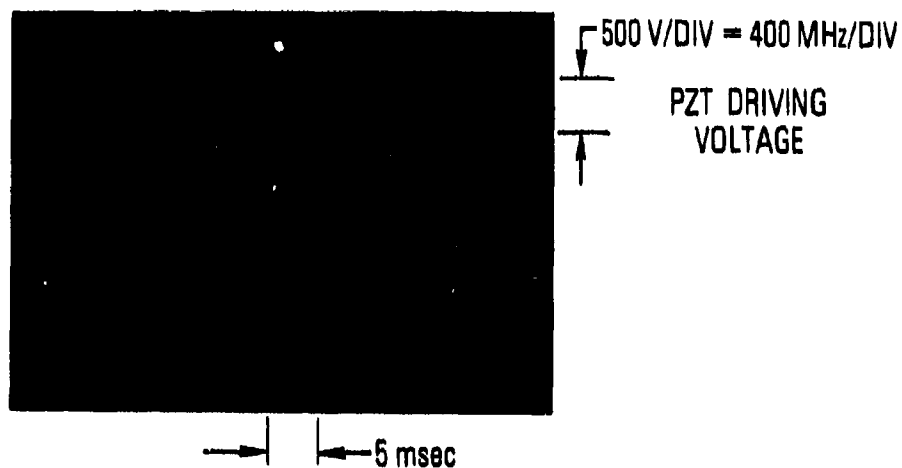


Figure 12. Frequency Scan of HF Laser Line  $P_2(8)$   
Showing Lamb Dip

## VI. CONCLUDING REMARKS

This small-scale subsonic flow laser has been used extensively.

Applications of this laser to date include:

1. The diagnostic gain scanning of a 2.5-kW cw HF(DF) chemical laser.<sup>2</sup>
2. Sun-laser heterodyne radiometer transatmospheric transmission measurement, HF(DF).<sup>3</sup>
3. Absorption diagnostic probing for electron-beam pulse laser (HF).
4. The study of window-material absorption.
5. Maritime atmospheric-transmission measurement.

It is expected that the laser will be used in HF(DF) laser kinetic rate studies. Most of these investigations depend principally upon the single-line amplitude and frequency stability characteristics of the laser and generally require  $< 100$  mW single-line power levels for adequate testing capability. Window material and optics coating studies, however, require higher power operation, and multiline performance is generally acceptable. In the event that power levels in excess of  $\sim 7$  to 15 W are required, side-by-side operation of two (or more) laser body channels can be used in order to double (or more) the gain length of the flow contained in the laser cavity.

The nominal order of magnitude increase in output power in HF, the even greater increase in DF laser power, and the increased number of observable lines of this laser relative to Hinchey's<sup>5</sup> earlier laser appear to be derived principally from the use of increased mass flows and the power supply, and the vacuum pumping capability of the present laser.

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